

# ELECTRONICS DEVICES AND CIRCUITS

## Section A

### Conducting Materials

## **OBJECTIVE**

**MOBILITY AND  
VARIATION OF  
MOBILITY WITH  
TEMPERATURE**

# ❖ Carrier Mobility

Macroscopic understanding

$$\mu = \frac{V_d}{E}$$

In a perfect Crystal

$$\rho = 0$$

$$\sigma \rightarrow \infty$$

It is a superconductor

Microscopic understanding? (what the carriers themselves are doing?)

$$\mu = \frac{q\tau}{m^*}$$

$$m_e^* < m_h^* \text{ in general}$$

$$m_e^*; n\text{-type}$$

$$m_h^*; p\text{-type}$$

# $\mu$

- A *perfect crystal has a perfect periodicity* and therefore the **potential** seen by a carrier in a perfect crystal is **completely periodic**.
- So the *crystal has no resistance to current flow* and behaves as a *superconductor*. The perfect periodic potential does not *impede the movement of the charge carriers*.

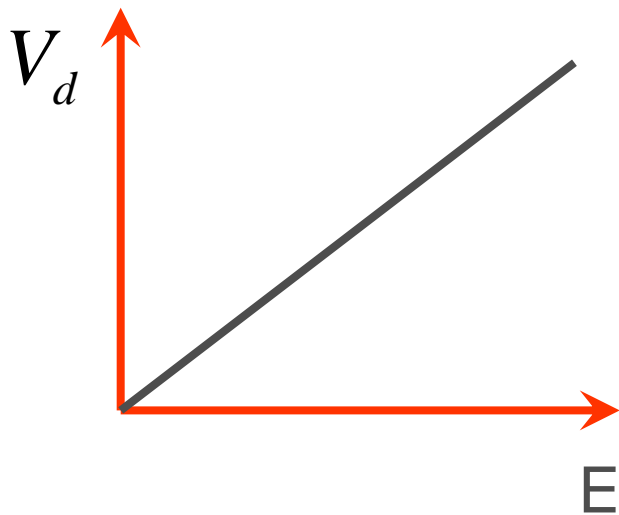
$\mu \dots$

- However, in a real device or specimen, the presence of impurities, interstitials, temperature, etc. creates a resistance to current flow.
- The presence of all these *upsets the periodicity of the potential* seen by a charge carrier.

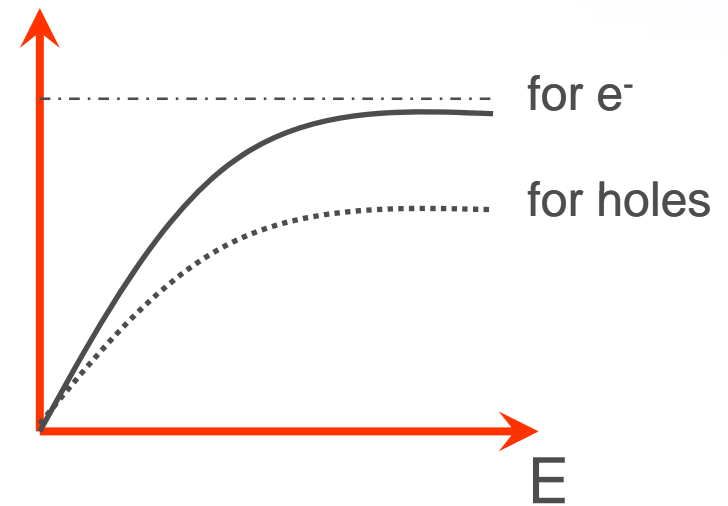
## ❖ Saturated Drift Velocities

$$V_d = \mu E$$

So one can make a carrier go as fast as we like just by increasing the electric field!!!



(a) Implication of above eqn.

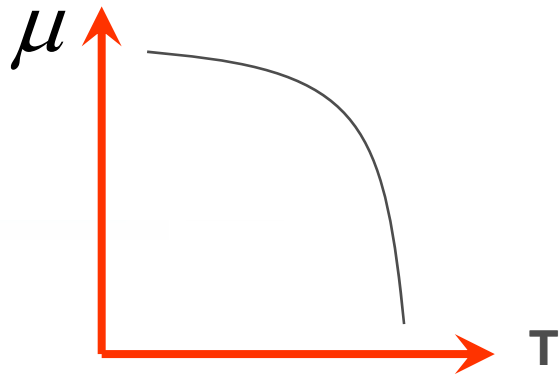


(b) Saturation drift velocity

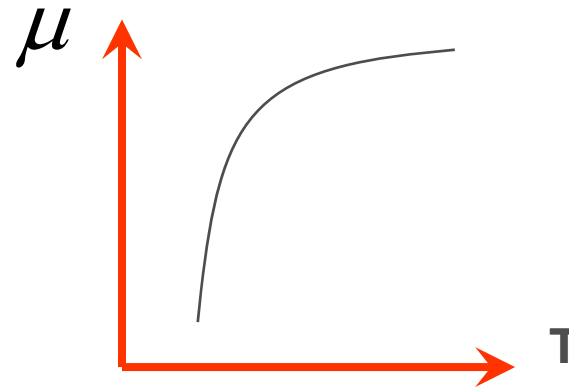
## ❖ Saturated Drift Velocities

- The equation of  $V_d = \mu \cdot E$  does not imply that  $V_d$  increases linearly with applied field  $E$ .
- $V_d$  increases linearly for low values of  $E$  and then it saturates at some value of  $V_d$  which is close  $V_{th}$  at higher values of  $E$ .
- Any further increase in  $E$  after saturation point does not increase  $V_d$  instead warms up the crystal.

## ❖ Mobility variation with temperature



High temperature

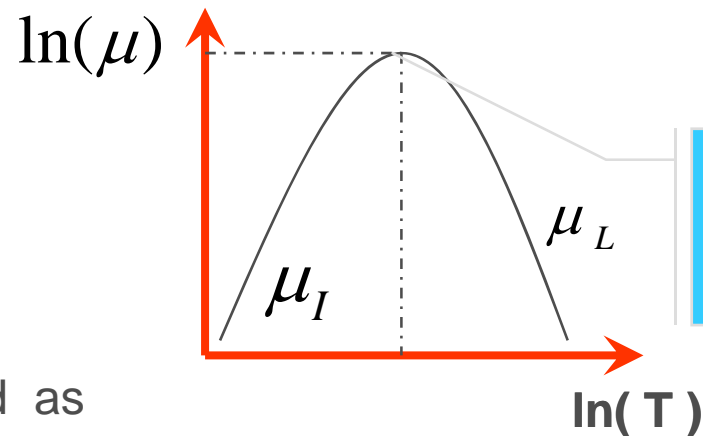


Low temperature

$$\frac{1}{\mu_T} = \frac{1}{\mu_L} + \frac{1}{\mu_I}$$



This equation is called as Mattheisen's rule.



Peak depends on the density of impurities